

Simultaneous single shot rotation-vibration non-equilibrium thermometry using pure rotational fs/ps CARS coherence beating

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We report the development of a simple and sensitive two beam hybrid femtosecond/picosecond pure rotational coherent anti-Stokes Raman scattering (fs/ps CARS) method to simultaneously measure the rotational and vibrational temperatures of diatomic molecules. Rotation-vibration non-equilibrium plays a key role in the chemistry and thermalization in low-temperature plasmas as well as thermal loading of hypersonic vehicles. This approach uses time-domain interferences between ground state and vibrationally excited N₂ molecules to intentionally induce coherence beating that leads to apparent non-Boltzmann distributions in the pure rotational spectra. These distortions enable simultaneous inference of both the rotational and vibrational temperatures. Coherence beating effects were observed in single shot fs/ps CARS measurements of a 75 Torr N₂ DC glow discharge and were successfully modeled for rotational and vibrational temperature extraction. We show that this method can be more sensitive than a pure rotational fs/ps CARS approach using a spectrally narrow probe pulse. Lastly, we experimentally measured the beat frequencies via Fourier transform of the time-domain response and obtained excellent agreement with the model. © 2022 Optical Society of America

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Rotation-vibration (R-V) non-equilibrium measurements are critical for understanding energy transfer and thermalization in low-temperature plasmas and hypersonic vehicles. In plasmas, energy transfer to molecules via collisions with electrons can induce significant vibrational non-equilibrium. Vibrationally-excited molecules are hypothesized to enhance chemical reactivity in plasma-catalysis [1], plasma-assisted combustion [2], and plasma CO₂ dissociation [3]. Furthermore, in hypersonic flows, shock-boundary layer interactions can induce R-V non-

equilibrium which affects the thermal loading and aerodynamics of hypersonic vehicles [4, 5]. Simultaneous measurements of both the rotational and vibrational temperatures are necessary to further our understanding in these areas.

Hybrid femtosecond/picosecond coherent anti-Stokes Raman scattering (fs/ps CARS) is a powerful laser diagnostic capable of probing gas temperature and species concentrations on the picosecond time scale with tens of microns of spatial resolution [6–17]. A method for retrieving vibrational and rotational temperatures simultaneously from spatially-resolved 1-D rotational fs/ps CARS images was recently developed [11]. This method uses the red-shift of rotational energy levels of vibrationally-excited molecules to measure the vibrational temperature along with the rotational temperature. The key advantages of this approach were the two-beam configuration for straightforward 1-D imaging and the elimination of the need for a separate Stokes pulse generated from an optical parametric amplifier (OPA). However, in [11] a 20 Hz repetition rate picosecond laser with a narrow line-width ($\sim 0.23 \text{ cm}^{-1}$) was used to resolve the one-sided broadening induced by vibrationally-excited N₂. It would be desirable generate the probe pulse from the femtosecond laser and enable kHz data acquisition rates without a separate laser.

Narrow probe picosecond pulses for fs/ps CARS have been previously generated from femtosecond pulses using second harmonic bandwidth compression (SHBC) [10, 16, 18, 19]. However, these probe pulses typically have 3 to 6 cm^{-1} bandwidths, while the energy spacing between adjacent vibrational levels in the pure rotational spectrum is approximately 1 cm^{-1} in N₂. Therefore, multiple Raman transitions from different vibrational levels will be coherently sampled within the probe pulse bandwidth and they can interfere with one another. This induces so-called “coherence beating” between these Raman transitions and distort the CARS spectrum as a function of probe delay. For example, in the rotational O₂ spectrum, coherence beating occurs due to the electronic triplet configuration of ground state O₂ which splits each rotational energy level into three closely spaced energy levels [20]. For certain probe delays, significant distortion of the CARS spectrum was observed, and accurate temperature

evaluations were not possible without considering the electronic triplet state splitting. In time-resolved CARS, frequency-domain spectra and Raman transition spacings have been obtained by taking the Fourier transform of beat patterns observed in probe delay scans on femtosecond to picosecond time scales [21, 22]. To mitigate coherence beating effects, dual-pump fs/ps CARS was used for R-V non-equilibrium measurements in [10, 23] and recently, simultaneous R-V non-equilibrium measurements were performed using fs/ps CARS in Mach 18 flows [17]. However, such measurements require multiple cameras and an optical parametric amplifier (OPA) to access the vibrational CARS spectrum. This adds complexity and reduces the pulse energy available for generating a CARS signal due to losses in the OPA. It would be desirable to use two-beam fs/ps CARS for R-V non-equilibrium measurements like in [11].

In this Letter, we present a new diagnostic method which instead exploits coherence beating to simultaneously infer the rotational and vibrational temperature from single shot pure rotational fs/ps CARS spectra. We simultaneously retain two-beam fs/ps CARS operation like in [11] and the potential for kHz data acquisition rates like in [15–17]. Furthermore, this approach is applicable to most diatomic molecules. We compare this approach to the frequency-domain one-sided broadening approach previously developed in [11] and show that the coherence beating can be even more sensitive than the one-sided broadening method. A 75 Torr N₂ DC glow discharge was used as a testbed for experimentally demonstrating this technique and vibrational temperatures ranging from 1000 K to 3600 K were measured on a single shot basis.

The detailed theory and principles behind hybrid fs/ps CARS frequency and time-domain modelling can be found in [9–12] and only a brief overview will be given here. Hybrid fs/ps CARS takes advantage of the broadband nature of femtosecond laser pulses to probe multiple Raman transitions simultaneously, while the non-resonant background is avoided with a time delayed spectrally-narrow probe pulse. In this study, two-beam phase matching is used where the both the pump and Stokes photons come from a single fs laser source. The fs/ps CARS time-domain model from [11] was used to calculate the spectral library for fitting the experimental spectra. The principle behind detecting vibrational and rotational energy distributions and their characteristic temperatures simultaneously from rotational CARS is given by Eq. 1 and 2:

$$F(v, J) = B_v(J(J+1)) - D_v(J^2(J+1)^2) \quad (1)$$

$$B_v = B_e - \alpha(v+1/2) + \gamma(v+1/2)^2 \quad (2)$$

where F is the rotational energy of an N₂ molecule in a state (v, J) with v and J being, respectively, the vibrational and rotational energy levels, B_v and D_v the vibrational energy level-dependent rotational and centrifugal constants, B_e is the equilibrium rotational constant, and α and γ are R-V coupling constants. From Eq. 1 and 2, the rotational energy levels exhibit a red-shift with increasing vibrational energy primarily due to the decrease in B_v . This vibrationally-induced shift can be on the order of 1 cm⁻¹. For molecules in R-V non-equilibrium, molecules in the excited vibrational states will appear at these shifted wavelengths. Such shifts have been observed before using spontaneous Raman scattering and nanosecond pulse-width CARS [24–26]. The coherence beating induced by interference between the neighboring vibrational levels allows for inference of the vibrational temperature in the pure rotational CARS spectrum [10, 11]. This approach is similar to [27], where an optimal probe delay for

vibrational Q-branch thermometry of N₂ was found due to the interferences between neighboring rotational lines.

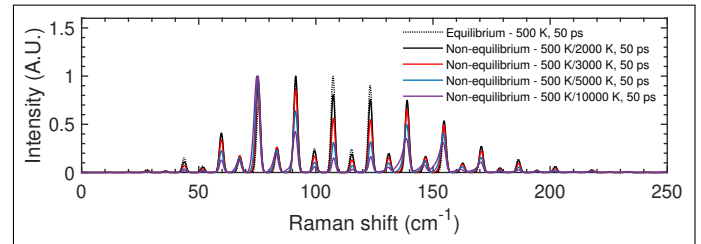


Fig. 1. Comparison of spectra in equilibrium and varying degrees of R-V non-equilibrium at a fixed probe delay of 50 ps.

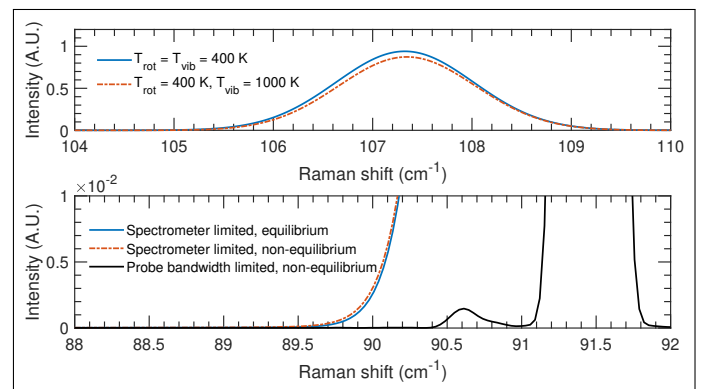


Fig. 2. Sensitivity comparison between coherence beating and one-sided broadening methods for non-equilibrium thermometry at $T_{rot} = 400$ K and $T_{vib} = 1000$ K.

Shown in Fig. 1 is a comparison of modelled hybrid fs/ps CARS spectra using a 10 picosecond pulse at an equilibrium temperature of 500 K and non-equilibrium vibrational temperatures up to 10,000 K. A stark difference in both profiles can be seen. The equilibrium spectrum produces an apparent Boltzmann distribution of rotational states. However, the non-equilibrium spectrum is strongly distorted, due to the presence of coherence beating between vibrationally excited N₂ and the ground state N₂. This distortion is a function of the vibrational temperature since the interference will be stronger with larger population of molecules in vibrationally excited states. This can be seen in Fig. 1 where the effects of coherence beating become more severe with higher vibrational temperatures. This suggests that this effect can be leveraged as a sensitive and quantitative diagnostic for simultaneous rotational and vibrational temperature measurements. Furthermore, these intensity distortions are much easier to detect than the one-sided broadening observed in [11]. Therefore, using the time-domain interferences can potentially be more sensitive than when using frequency domain detection. In Fig. 2, we compare the two approaches at a 50 ps delay and a vibrational temperature of 1000 K. If the frequency resolution of the measurement is limited by the probe bandwidth and not the spectrometer, then both techniques can measure this condition. However, typical grating spectrometers have instrument spectral broadening functions of 0.5 to 1 cm⁻¹. In this case, the instrument broadening smears the adjacent peaks and limits the sensitivity of the measurement as shown in the Fig. 2. The equilibrium and non-equilibrium instrument broadened spectra are nearly indistinguishable at this vibrational temperature.

142 For coherence beating, this is not a problem, and the theoretical
143 sensitivity is retained.

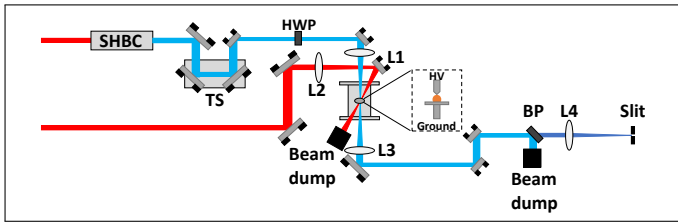


Fig. 3. Experimental fs/ps CARS schematic. L1 to L3: $f = 400$ mm, L4: $f = 75$ mm, BP: angle-tuned bandpass filter, TS: translation stage, HWP: half-wave plate.

144 To experimentally demonstrate this technique, we measured
145 fs/ps CARS spectra on a single shot basis in a 75 Torr N_2 DC
146 glow discharge plasma. Due to the low reduced electric field,
147 efficient pumping of vibrational energy levels induced R-V non-
148 equilibrium in this plasma. However, the DC glow discharge
149 drifted in space, so single shot CARS measurements were neces-
150 sary since averaging was not possible. The hybrid fs/ps CARS
151 setup is shown in Fig. 3. The 80 fs pump/Stokes pulse centered
152 around 800 nm were supplied by a femtosecond Ti:Sapphire re-
153 generative amplifier (Astrella, Coherent), and a 9 ps probe pulse
154 centered around 399.6 nm was generated using a SHBC design
155 similar to [28]. The probe pulse shape was measured using
156 cross-correlation in pure CH_4 and assumed to be transform-
157 limited (1.63 cm^{-1}). A 1500 gr/mm diffraction grating with
158 90% diffraction efficiency (Spectrogon) was used as the disper-
159 sive element to stretch the femtosecond pulse and the stretched
160 beam was split via a 50/50 beam splitter (Thorlabs). Propaga-
161 tion through the beam splitter and reflection by the beam splitter
162 produces two spatially chirped pulses that are mirror images
163 of each other. In other words, they are oppositely chirped, as
164 required for spectral compression via SHBC. A 250 mm focal
165 length achromatic lens was used to collimate the dispersed light.
166 Since the dispersion was only in one axis, a sheet was formed at
167 the Fourier plane, where the beta barium borate (BBO, Eksma
168 Optics) crystal was placed. Temporal alignment of the split
169 pulses was achieved through adjusting a micrometer transla-
170 tion stage and probe pulses with energies up to $100 \mu\text{J}$ were
171 generated. The femtosecond pulse and the picosecond pulse
172 were spherically focused separately and directed to intersect at
173 a 3-degree angle in a two-beam phase matching configuration.
174 The spatial overlap was approximately 1 mm. The time delay
175 between the probe pulse and pump/Stokes pulse was set with a
176 delay line on a motorized translation stage (LTS150, Thorlabs).
177 Detection of the CARS signal was performed via two angle-
178 tuned bandpass filters (Semrock), a Czerny-Turner spectrometer
179 (SP2500i, Princeton Instruments), and an intensified charge cou-
180 pled device (ICCD, PIMAX1300, Princeton Instruments). The
181 instrument function of the spectrometer was determined via
182 fitting room-temperature N_2 spectra. The plasma reactor was
183 the same as in [11], but the electrodes were swapped for a pin
184 to plane discharge geometry with a pin anode and a 1 cm di-
185 ameter cylindrical cathode separated by 7 mm. The anode was
186 connected to a high voltage direct current (DC) power supply
187 (Glassman) in series with a 100 kOhm ballast resistor. The volt-
188 age at the anode was 500 V and the potential drop across the
189 resistor was 1500 V. The synchronization between the laser and
190 the ICCD was achieved using a delay generator (DG645, Stan-

ford Research Systems) triggered by the femtosecond laser. A
192 constant flow of 600 SCCM of pure N_2 was maintained using a
193 mass flow controller (MKS) and the pressure was kept constant
194 with a downstream automatic throttle valve (153D, MKS).

195 Shown in Fig. 4 are the experimental spectra of this glow
196 discharge at different probe pulse delays as well as the non-
197 equilibrium fits. The effects of coherence beating on the spectra
198 are immediately obvious from comparison of the data and the
199 R-V equilibrium simulation. Varying the probe delays induced
200 different beat patterns in the rotational spectrum and these co-
201 herence beating effects were successfully fit by the CARS model.
202 For the 92 ps delay, a clear deviation from a Boltzmann distribu-
203 tion of rotational states can be seen, while at 44 ps, the distor-
204 tion is less obvious. This suggests that, like in [27], there should be
205 an optimal probe delay to maximize measurement sensitivity.

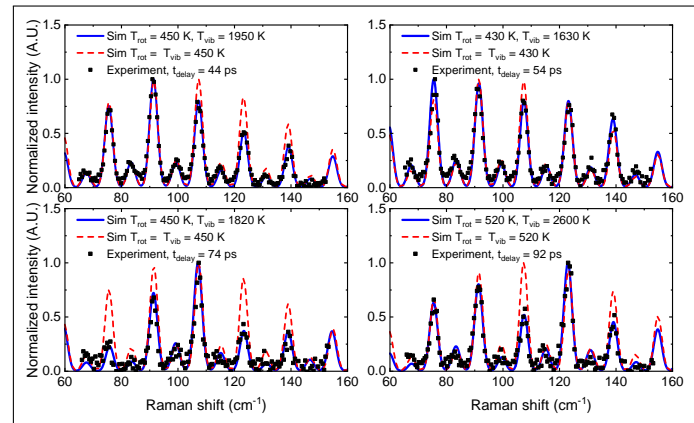


Fig. 4. Single shot fs/ps CARS spectra of a N_2 glow discharge and corresponding simulated fits for R-V non-equilibrium at different probe delays. Equilibrium spectra at the same rotational temperatures were plotted for reference.

206 The time-dependent R-V non-equilibrium using probe delays
207 of 74 ps and 92 ps were measured in the discharge as shown in
208 Fig. 5. Only spectra with signal to noise ratios above 25 were
209 accepted for temperature evaluation and the 95% confidence
210 intervals of the fits were used as error bars. Both acquisition
211 runs show similar degrees of R-V non-equilibrium and tempo-
212 ral dynamics. The rotational temperature mostly ranged from 350
213 to 500 K and the vibrational temperature was measured to be
214 within 1000 to 1500 K. However, in both time traces, brief peaks
215 in vibrational temperature up to 3600 K were observed. This
216 likely was due to optimal overlap between the glow discharge
217 and the CARS probe volume.

218 Lastly, we scanned the probe delay across 300 ps with 10
219 ps steps to perform time-resolved CARS and used the Fourier
220 transform to retrieve the beat frequencies. Plotted in Fig. 6 are
221 the measured and theoretical beat frequencies. Differences of
222 less than 0.2 cm^{-1} between the theory and experiment was ob-
223 tained. Considering that the frequency resolution of the Fourier
224 transform was $\sim 0.1 \text{ cm}^{-1}$, this is excellent agreement.

225 For the present approach, there are several challenges and op-
226 portunities for further exploration. Due to the probe bandwidth,
227 individual hot bands remain unresolved, and non-Boltzmann
228 distributions may affect the observed spectra. With the current
229 probe pulse, only vibrational levels up to $v=2$ are within the
230 probe bandwidth. Therefore, these measurements are primarily
231 sensitive to the first-level vibrational temperature. Additionally,

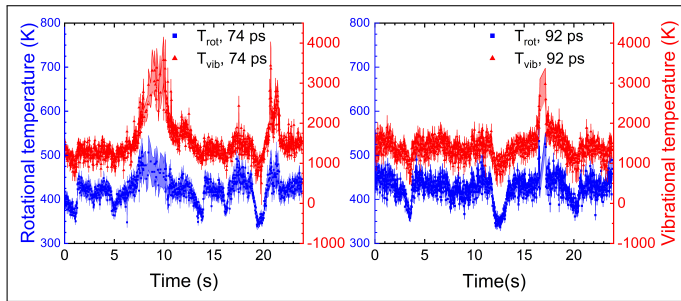


Fig. 5. Time dependent R-V non-equilibrium measured in a N_2 DC glow discharge with probe delays of 74 ps (left) and 92 ps (right). Error bars are shaded.

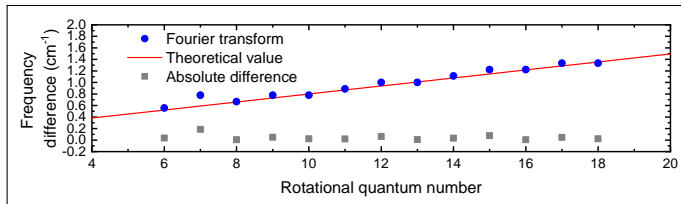


Fig. 6. Measured beat frequencies via Fourier transform of the probe delay scan compared to those calculated by Eq. 1 and 2.

measurements in environments above atmospheric pressure may suffer from signal loss from collisional dephasing prior to the first beat period. Finally, spectral overlap with other molecular species could mask the coherence beating. However, not all lines will overlap, and such effects can be mitigated with a probe delay aligned to the beat periods of spectrally isolated rotational transitions. This may also present an opportunity to simultaneously measure the coherence beating and R-V non-equilibrium of multiple species with an optimized probe delay.

In summary, a new method to simultaneously measure rotational and vibrational temperatures of molecules using coherence beating of pure rotational hybrid fs/ps CARS was developed. The distortion of the pure rotational fs/ps CARS spectrum increased with the vibrational temperature, which indicated that both rotational and vibrational temperature can be directly retrieved from the pure rotational fs/ps CARS spectrum. Coherence beating between rotational lines of vibrationally excited and ground state N_2 was observed experimentally in a N_2 DC glow discharge and the measured beat patterns were in excellent agreement with the model. Furthermore, the extraction of the rotational and vibrational temperatures from fitting the fs/ps CARS spectra distorted by the coherence beating was demonstrated. This technique was shown to be more sensitive than using one-sided broadening like in [11]. Moreover, the switch to measurement of intensity distortions in the spectra enables the possibility of sensitive 2-D R-V non-equilibrium measurements using hyperspectral 2D-CARS [8, 13] as well as high speed kHz two-beam fs/ps CARS of R-V non-equilibrium in plasmas and hypersonic wind tunnels.

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Data availability. Data underlying the results presented in this paper are not publicly available at this time but may be obtained from the authors upon reasonable request.

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